

Explicit formula for singular vectors of the Virasoro algebra with central charge less than 1

Reiho Sakamoto

*Department of Physics, Graduate School of Science
University of Tokyo, Hongo, Bunkyo-ku, Tokyo, 113-0022, Japan*

reiho@monet.phys.s.u-tokyo.ac.jp

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Abstract

We calculate explicitly the singular vectors of the Virasoro algebra with the central charge $c \leq 1$. As a result, we have an infinite sequence of the singular vectors for each Fock space with given central charge and highest weight, and all its elements can be written in terms of the Jack symmetric functions with rectangular Young diagram.

Since the paper of Belavin, Polyakov and Zamolodchikov [1], singular vectors of the Virasoro algebra attract a lot of attention because of its relationship with correlation functions. Many attempts were done to make the explicit expression of these vectors in the Verma module (see for example [2-4]), however it remains to be unclear what is the structure of these vectors. On the other hand, singular vectors on the Fock space are known to have relationships with symmetric functions. Goldstone, Wakimoto and H. Yamada showed that when the central charge c is equal to 1, then the singular vectors of smallest degree are proportional to the Schur symmetric functions with rectangular Young diagram [5, 6]. Using the actions of the Virasoro generators on the Schur symmetric

functions, [7] generalized this result to the case $c < 1$ and expressed the singular vectors as a sum of the Schur symmetric functions. Mimachi and Y. Yamada extended these results to the general value of c and showed that the singular vector of the Fock space $\mathcal{F}_{A_{r+1}, s+1}$ (see below for definition) with degree rs is proportional to the Jack symmetric function with rectangular Young diagram (s^r) [8, 9]. However, as is well known, there are infinitely many singular vectors in each Fock space $\mathcal{F}_{A_{r+1}, s+1}$ if $c \leq 1$, and all above results determine only the singular vectors of smallest degree. Recently, [10] demonstrated that the set of all the bosonized Jack symmetric functions forms a natural basis of the Fock space, and conjectured explicitly the actions of the Virasoro generators on this Jack basis. In this paper, applying these results, we determine the explicit expression of all the singular vectors of the Fock space $\mathcal{F}_{A_{r+1}, s+1}$ when central charge is less than 1. As a result, we found that all the singular vectors are proportional to the Jack symmetric functions with rectangular Young diagram, which indicates further deep relationships between the Jack symmetric functions and the representations of the Virasoro algebra.

The set-up is as follows. Denote the Virasoro generators as usual by L_n ($n \in \mathbf{Z}$) and the central charge as c , with commutation relations

$$[L_n, L_m] = (n - m)L_{n+m} + \frac{n(n^2 - 1)}{12} c \delta_{n+m, 0}. \quad (1)$$

As a representation space, we take the Fock space \mathcal{F}_A defined by

$$\mathcal{F}_A := \mathbf{C}[a_{-1}, a_{-2}, a_{-3}, \dots] |A\rangle, \quad (2)$$

where we use the bosonic operators $[a_n, a_m] = n \delta_{n+m, 0}$ ($n, m \in \mathbf{Z}$), and a vacuum vector $|A\rangle$ defined by

$$a_0 |A\rangle = A |A\rangle, \quad a_n |A\rangle = 0 \quad (n \in \mathbf{Z}_{>0}) \quad (3)$$

Then the bosonic representation of L_n 's is given as follows — the Feigin-Fuchs representation [11],

$$L_n = \frac{1}{2} \sum_{k \in \mathbf{Z}} : a_{n-k} a_k : - A_{1,1} (n+1) a_n, \quad (4)$$

here we use the notation

$$A_{r,s} = \frac{1}{\sqrt{2}} \left(r\sqrt{\beta} - \frac{s}{\sqrt{\beta}} \right). \quad (5)$$

This β is related to the central charge c as

$$c = 1 - \frac{6(1-\beta)^2}{\beta}. \quad (6)$$

Note that by the universality of the Verma module, we have a map from the Verma module with highest weight

$$h_{r,s} = \frac{(r\beta - s)^2 - (\beta - 1)^2}{4\beta}, \quad (7)$$

to the Fock space $\mathcal{F}_{A_{r+1,s+1}}$ (we have $L_0|A_{r+1,s+1}\rangle = h_{r,s}|A_{r+1,s+1}\rangle$).

Another ingredient of this study is the Jack symmetric function. Consider the Hamiltonian

$$H_\beta := \sum_{i=1}^N \left(x_i \frac{\partial}{\partial x_i} \right)^2 + \beta \sum_{i < j} \frac{x_i + x_j}{x_i - x_j} \left(x_i \frac{\partial}{\partial x_i} - x_j \frac{\partial}{\partial x_j} \right), \quad (8)$$

where N is a number of variables. Then the Schrödinger equation

$$H_\beta J_\lambda(x_1, x_2, \dots, x_N) = \epsilon_\lambda J_\lambda(x_1, x_2, \dots, x_N), \quad (9)$$

$$\epsilon_\lambda = \sum_{i=1}^N \left(\lambda_i^2 + \beta(N+1-2i)\lambda_i \right), \quad (10)$$

is known to be exactly solvable and its wave functions are called the Jack symmetric polynomials [12, 13]. These J_λ 's are the N variable symmetric polynomials parameterized by partitions $\lambda = (\lambda_1, \lambda_2, \lambda_3, \dots)$. Physically, Jack polynomials express the excited states of the Calogero-Sutherland model [14-17] with periodic boundary condition. Ujino, Hikami and Wadati investigated the complete quantum integrability of these systems [18, 19] (see also [20]).

Instead of using the coordinates x_i , we can express symmetric polynomials uniquely in terms of power sums $p_n := \sum x_i^n$, if the number of variables is sufficiently large. After expressing the Jack polynomials by power sums, we can take the number of variables as infinity. These Jack symmetric polynomials with infinitely many variables are called Jack symmetric functions.

To characterize Jack symmetric functions by orthogonality, we define the inner product on the space of symmetric functions as

$$\langle p_\lambda, p_\mu \rangle := \delta_{\lambda, \mu} z_\lambda \beta^{-l(\lambda)}, \quad (11)$$

$$p_\lambda := \prod_{i=1}^{l(\lambda)} p_{\lambda_i}, \quad (12)$$

$$z_\lambda := \prod_i i^{m_i} \cdot m_i!. \quad (13)$$

Here, $l(\lambda)$ is a length of partition λ and m_i 's are the multiplicity of i in partition λ , i.e. we have $\lambda = (\cdots 3^{m_3}, 2^{m_2}, 1^{m_1})$. We also denote the weight of partition λ as $|\lambda|$ and define the coordinate of partition like a matrix, i.e. if the box is on row i and column j , then the coordinate is (i, j) . Using this coordinate, define the arm length $a(s)$ and leg length $l(s)$ of the box $s = (i, j) \in \lambda$ by

$$a(s) := \lambda_i - j, \quad l(s) := \lambda'_j - i, \quad (14)$$

where λ'_j is the j -th element of the transposed partition λ' .

Then our normalization J_λ is characterized by its norm [13]

$$\langle J_\lambda, J_\mu \rangle = \delta_{\lambda, \mu} \frac{1}{\beta^{2|\lambda|}} \prod_{s \in \lambda} (a(s) + (l(s) + 1)\beta) ((a(s) + 1) + l(s)\beta), \quad (15)$$

especially we set the coefficient in front of $p_1^{|\lambda|}$ to be unity. See [21] for rigorous discussion about uniqueness. As an example, we give a list of Jack symmetric function with degree 3 (here we set $\alpha := 1/\beta$),

$$J_{\square\square\square} = 2\alpha^2 p_3 + 3\alpha p_2 p_1 + p_1^3,$$

$$J_{\square\square} = -\alpha p_3 + (\alpha - 1)p_2 p_1 + p_1^3,$$

$$J_{\square} = 2p_3 - 3p_2 p_1 + p_1^3.$$

From this characterization, we can identify the Jack symmetric function and element of the Fock space [22, 23]. To do this, we use the following two identifications,

$$p_n \longleftrightarrow \sqrt{\frac{2}{\beta}} a_{-n} |a\rangle, \quad (16)$$

$$p_n \longleftrightarrow \langle A | a_n \frac{1}{\sqrt{2\beta}}. \quad (17)$$

In [10], it was proposed that we should take the basis of the Fock space as a set of all the bosonized Jack symmetric functions,

$$\mathcal{F}_A = \text{span}\{J_\lambda|A\rangle \mid \text{any partition } \lambda\}. \quad (18)$$

Since [10] deals with only the leftwards action of L_n operators on this Jack basis, we need to rewrite the formula for rightwards action. In subsequent discussion, we write $J_\lambda|A_{r+1,s+1}\rangle = |J_\lambda\rangle$ for simplicity, and also denote the action of L_n as

$$\langle J_\lambda|L_n = \sum_{|\mu| - |\lambda| = n} \langle J_\mu|c_{\lambda\mu}, \quad (19)$$

$$L_n|J_\mu\rangle = \sum_{|\mu| - |\lambda| = n} d_{\mu\lambda}|J_\lambda\rangle. \quad (20)$$

Note that $c_{\lambda\mu}$'s can be determined from the formula in [10]. Then from the identity

$$\begin{aligned} (\langle J_\lambda|L_n)|J_\mu\rangle &= c_{\lambda\mu}\langle J_\mu|J_\mu\rangle \\ &= \langle J_\lambda|(L_n|J_\mu\rangle) = d_{\mu\lambda}\langle J_\lambda|J_\lambda\rangle, \end{aligned} \quad (21)$$

we have the following relation,

$$d_{\mu\lambda} = \frac{\langle J_\mu|J_\mu\rangle}{\langle J_\lambda|J_\lambda\rangle} c_{\lambda\mu}. \quad (22)$$

First we consider the action of L_1 on Jack symmetric function with rectangular partition. In this case, we have

$$\begin{aligned} c_{(n^{m-1}, n-1), (n^m)} &= c_{\begin{array}{|c|c|} \hline \square & \square \\ \hline \end{array}} \\ &= \sqrt{2\beta} \frac{1 \cdot \beta}{(1 + (m-1)\beta)((n-1) + \beta)} A_{r-m, s-n}, \end{aligned} \quad (23)$$

$$\frac{\langle J_{(n^m)}|J_{(n^m)}\rangle}{\langle J_{(n^{m-1}, n-1)}|J_{(n^{m-1}, n-1)}\rangle} = \frac{1}{\beta^2} \cdot \frac{(0 + m\beta)(1 + (m-1)\beta)((n-1) + \beta)(n + 0\beta)}{(0 + \beta)(1 + 0\beta)}. \quad (24)$$

Note that the leftward action $\langle J_\lambda|L_1$ is amount to the addition of one box to the partition λ . Then from the orthogonality of the Jack symmetric functions, we conclude that the rightward action $L_1|J_\lambda\rangle$ amount to subtract one box from the partition λ . Thus we obtain the formula for $J_{(n^m)}$ as

$$\begin{aligned} L_1|J_{(n^m)}\rangle &= c_{(n^{m-1}, n-1), (n^m)} \frac{\langle J_{(n^m)}|J_{(n^m)}\rangle}{\langle J_{(n^{m-1}, n-1)}|J_{(n^{m-1}, n-1)}\rangle} |J_{(n^{m-1}, n-1)}\rangle \\ &= \sqrt{\frac{2}{\beta}} \cdot \frac{1}{\beta} mn\beta A_{r-m, s-n} |J_{(n^{m-1}, n-1)}\rangle. \end{aligned} \quad (25)$$

Next we consider the action of L_2 . In this case, we have the following equalities,

$$\begin{aligned}
c_{(n^{m-1}, n-2), (n^m)} &= c_{\begin{array}{|c|c|} \hline \square & \square \\ \hline \end{array}} \\
&= \sqrt{2\beta} \beta \frac{2\beta}{(2 + (m-1)\beta)((n-2) + \beta)} \\
&\quad \times \frac{1}{(1 + (m-1)\beta)((n-1) + \beta)} A_{r-m, s-n}, \tag{26}
\end{aligned}$$

$$\begin{aligned}
\frac{\langle J_{(n^m)} | J_{(n^m)} \rangle}{\langle J_{(n^{m-1}, n-2)} | J_{(n^{m-1}, n-2)} \rangle} &= \frac{1}{\beta^4} \cdot \frac{1}{(1 + \beta)(2 + 0\beta)(0 + \beta)(1 + 0\beta)} \\
&\quad \times (1 + m\beta)(2 + (m-1)\beta)(0 + m\beta)(1 + (m-1)\beta) \\
&\quad \times ((n-1) + \beta)(n + 0\beta)((n-2) + \beta)((n-1) + 0\beta), \tag{27}
\end{aligned}$$

$$\begin{aligned}
c_{(n^{m-2}, (n-1)^2), (n^m)} &= c_{\begin{array}{|c|c|} \hline \square & \square \\ \hline \end{array}} \\
&= -\sqrt{2\beta} \frac{1 \cdot 2\beta}{(1 + (m-2)\beta)((n-1) + 2\beta)} \\
&\quad \times \frac{\beta}{(1 + (m-1)\beta)((n-1) + \beta)} A_{r-m, s-n}, \tag{28}
\end{aligned}$$

$$\begin{aligned}
\frac{\langle J_{(n^m)} | J_{(n^m)} \rangle}{\langle J_{(n^{m-2}, (n-1)^2)} | J_{(n^{m-2}, (n-1)^2)} \rangle} &= \frac{1}{\beta^4} \cdot \frac{1}{(0 + 2\beta)(1 + \beta)(0 + \beta)(1 + 0\beta)} \\
&\quad \times (0 + m\beta)(1 + (m-1)\beta)(0 + (m-1)\beta)(1 + (m-2)\beta) \\
&\quad \times ((n-1) + 2\beta)(n + \beta)((n-1) + \beta)(n + 0\beta). \tag{29}
\end{aligned}$$

Combining these equalities, we have

$$\begin{aligned}
L_2 |J_{(n^m)}\rangle &= c_{(n^{m-1}, n-2), (n^m)} \frac{\langle J_{(n^m)} | J_{(n^m)} \rangle}{\langle J_{(n^{m-1}, n-2)} | J_{(n^{m-1}, n-2)} \rangle} |J_{(n^{m-1}, n-2)}\rangle \\
&\quad + c_{(n^{m-2}, (n-1)^2), (n^m)} \frac{\langle J_{(n^m)} | J_{(n^m)} \rangle}{\langle J_{(n^{m-2}, (n-1)^2)} | J_{(n^{m-2}, (n-1)^2)} \rangle} |J_{(n^{m-2}, (n-1)^2)}\rangle \\
&= \sqrt{\frac{2}{\beta}} \cdot \frac{1}{\beta^2} \left(mn\beta \cdot \frac{(n-1)(1 + m\beta)}{(1 + \beta)} A_{r-m, s-n} |J_{(n^{m-1}, n-2)}\rangle \right. \\
&\quad \left. - mn\beta \cdot \frac{(m-1)\beta(n + \beta)}{(1 + \beta)} A_{r-m, s-n} |J_{(n^{m-2}, (n-1)^2)}\rangle \right). \tag{30}
\end{aligned}$$

From now on, we restrict the value of β as follows,

$$\beta = \frac{p}{q} > 0, \tag{31}$$

where p and q are mutually prime positive integers. When we use this value of β , the central charge of the Virasoro algebra satisfies $c \leq 1$. If we rewrite the definition of $A_{r,s}$

in this case as

$$A_{r,s} = \frac{1}{\sqrt{2\beta}} \left(\frac{pr}{q} - s \right), \quad (32)$$

we have the equality

$$A_{m-(m+aq), n-(n+ap)} = A_{-aq, -ap} = 0, \quad a \in \mathbf{Z}_{>0}. \quad (33)$$

Combining this simple equality with equations (25) and (30), we obtain the following remarkable result.

Theorem *If $\beta = p/q > 0$, then the Fock space $\mathcal{F}_{A_{m+1}, n+1}$ ($m, n \in \mathbf{Z}_{>0}$) have the following infinite sequence of singular vectors;*

$$|J_{(n^m)}\rangle, |J_{(n+p)^{(m+q)}}\rangle, |J_{(n+2p)^{(m+2q)}}\rangle, \dots, |J_{(n+ap)^{(m+aq)}}\rangle, \dots \quad (34)$$

To verify this, we consider the action of L_1 and L_2 on $|J_{(n+ap)^{(m+aq)}}\rangle$. Then from equations (25) and (30), we have the term $A_{m-(m+aq), n-(n+ap)}$, which is equal to 0 by (33).

We remark that the first element of the sequence (34) coincide with the result given in [9] (see also [24, 25] for the structure of the Verma modules). We also remark that in (34), we do not take care of the precise normalization.

Example We consider the case $\beta = 3/4$. This amount to $c = 1/2$ to which the critical Ising model corresponds. If we choose the highest weight as $A_{1+1, 2+1}$, then we have a sequence of singular vectors

$$|J_{(2)}\rangle, |J_{(5^5)}\rangle, |J_{(8^9)}\rangle, |J_{(11^{13})}\rangle, \dots \quad (35)$$

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